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**Diffeomorphisms constructively associated with
mutually diverging spacetimes which allow a natural
identification of event points in general relativity.
Part II**

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Fisica matematica. — *Diffeomorphisms constructively associated with mutually diverging spacetimes which allow a natural identification of event points in general relativity.* Part II. Nota di GAETANO ZAMPIERI (*), presentata (**), dal Corrisp. A. BRESSAN.

RIASSUNTO. — In questo lavoro si dà una definizione di divergenza fra cronotopi della Relatività Generale e si costruisce un criterio per l'identificazione dei punti eventi di cronotopi divergenti che appartengono ad una classe consistente con la presenza di campi elettromagnetici nel vuoto.

4. INTRODUCTION ⁽¹⁾.

In Sect. 5 the diffeomorphisms that provide the quasi-absolute concept of event point are constructed. To be more detailed, let us consider the typical equivalence class \mathcal{D} of the divergence relation, assume that $S_4^1, S_4^2 \in \mathcal{D}$, and that P^{21} be a region of type \mathcal{D} , with respect to both spacetimes, where the structures of S_4^1 and S_4^2 coincide. Suppose $\mathcal{E}^1 \in S_4^1$ and consider the maximal integral line c_1 of the reference frame η^1 on S_4^1 , starting from \mathcal{E}^1 ⁽²⁾ (whose parameter is the proper time ⁽³⁾).

On the basis of the definition of a type \mathcal{D} region, there is a value $-\tau$ of the parameter, such that $\tau \geq 0$ and $\mathcal{E} = c_1(-\tau) \in P^{21}$.

Now consider the maximal integral line c_2 , of the reference frame η^2 on S_4^2 , starting from \mathcal{E} . If its parameter takes on the value τ , then the event $\mathcal{E}^2 = c_2(\tau)$ is by definition the correspondent of \mathcal{E}^1 in S_4^2 . Otherwise, \mathcal{E}^1 has no correspondent.

This definition makes sense because \mathcal{E}^2 , or its inexistence, is independent of the choice of P^{21} and τ .

Thus a bijection f^{21} , from a subset of S_4^1 onto a subset of S_4^2 , is defined. It turns out to be a diffeomorphism whose restriction to P^{21} is the inclusion map. Moreover, f^{21} is the identity map if the two spacetimes coincide, and $f^{21} = (f^{12})^{-1}$, where f^{12} is the diffeomorphism associated with the inverse

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(1) This section is the sequel of Sect. 1 in Part. I—see ref. [7].

(2) We have (i) $c_1 : \mathfrak{J}(\mathcal{E}^1) \rightarrow S_4^1$, where $\mathfrak{J}(\mathcal{E}^1)$ is an open interval of \mathbf{R} containing zero, (ii) $c_1(0) = \mathcal{E}^1$, (iii) $\dot{c}_1(\tau) = \eta^1[c_1(\tau)] \forall \tau \in \mathfrak{J}(\mathcal{E}^1)$, where $\dot{c}_1(\tau)$ is the tangent vector at $c_1(\tau)$, and (iv) c_1 is maximal—see e.g. ref [5] p. 4.

(3) See e.g. ref. [5] p. 41.

pair (S_4^2, S_4^1) . Finally, if \mathcal{E}^1 has a correspondent \mathcal{E}^2 , and \mathcal{E}^2 has a correspondent \mathcal{E}^3 in a third spacetime $S_4^3 \in \mathcal{D}$, then \mathcal{E}^3 corresponds to \mathcal{E}^1 under the diffeomorphism f^{31} associated with (S_4^1, S_4^3) ⁽⁴⁾.

The section ends with the definition of a second class of spacetimes, \mathfrak{B} , whose typical element S_4 belongs to \mathfrak{A} , and is such that the reference frame η is a complete vector field ⁽⁵⁾.

In this case the diffeomorphism f^{21} maps the first spacetime onto the second. Thus in the class \mathfrak{B} we have an absolute concept of event point.

5. DIFFEOMORPHISMS ASSOCIATED WITH MUTUALLY DIVERGING SPACETIMES

Let S_4^k ($k = 1, 2$) be spacetimes in \mathfrak{A} —see ref. [7] Def. 3—and mutually diverging—see ref. [7] Def. 2. Let P^{21} be as in ref. [7] Def. 2, and let η^k ($k = 1, 2$) be the analogue for S_4^k of the reference frame η on S_4 mentioned in [7] Def. 3.

LEMMA 2. *Under these assumptions*

$$\eta^1|_{P^{21}} = \eta^2|_{P^{21}}.$$

Proof. If $\mathcal{E} \in P^{21}$ and $k \in \{1, 2\}$, then, by Def. 1 in [7], the causal past of \mathcal{E} in S_4^k coincides with the one in P^{21} endowed with the structure induced from the structure of S_4^k ⁽⁶⁾. Furthermore, the structures of S_4^1 and S_4^2 coincide on P^{21} —see [7] Def. 2. This, and the definition of $\eta^k(\mathcal{E})$ —which depends only on the causal past of \mathcal{E} —, yield the proposition. q.e.d.

Let $\mathcal{E}^1 \in S_4^1$ and let $c_1: \mathfrak{J}(\mathcal{E}^1) \rightarrow S_4^1$ be the maximal integral line of η^1 starting from \mathcal{E}^1 (i.e. with $c_1(0) = \mathcal{E}^1$)—see fnt (2). Then the map $\bar{c}: \{\alpha \in \mathbf{R} : \alpha \geq 0, -\alpha \in \mathfrak{J}(\mathcal{E}^1)\} \rightarrow S_4^1, \alpha \mapsto c_1(-\alpha)$ is a causal past-pointing and past-inextendible curve ⁽⁷⁾; and it enters P^{21} and then remains within it because $P^{21} \in \mathcal{D}(S_4^1)$ —see Def. 1 in [7]. Therefore we can consider $\tau \in \mathbf{R}$ with $\tau \geq 0, -\tau \in \mathfrak{J}(\mathcal{E}^1)$, and $c_1(-\tau) \in P^{21}$. Furthermore,

$$(1) \quad \tau' \geq \tau, -\tau' \in \mathfrak{J}(\mathcal{E}^1) \Rightarrow c_1(-\tau') \in P^{21}.$$

(4) The event \mathcal{E}^1 can have a correspondent \mathcal{E}^3 in S_4^3 without any correspondent in S_4^2 .

(5) That is the domain of all its maximal integral curves is the whole \mathbf{R} .

(6) In fact the causal past—see [4] p. 183—of \mathcal{E} coincides with the union (of the ranges) of the past-pointing and past-inextendible causal curves starting from \mathcal{E} .

(7) \bar{c} is past-inextendible, i.e. it has no past-endpoint—see fnt 7 in Part I i.e. [7]—because it is a maximal integral curve of a causal vector field (which is nowhere vanishing).

Now let $\mathcal{E} = c_1(-\tau)$ and let $c_2 : \mathfrak{J}(\mathcal{E}) \rightarrow S_4^2$ be the maximal integral curve of the vector field η^2 starting from \mathcal{E} (i.e. with $c_2(0) = \mathcal{E}$).

If $\tau \in \mathfrak{J}(\mathcal{E})$, then I say that \mathcal{E}^1 is R^{21} -related to $\mathcal{E}^2 = c_2(\tau)$.

This definition is good because \mathcal{E}^2 , or its inexistence, is independent of the aforementioned choice of P^{21} and τ . In fact, let \bar{P}^{21} and $\bar{\tau}$ satisfy the same conditions, and let us suppose that $\bar{\tau} > \tau$. Then (1) yields $c_1(-\tau') \in P^{21}$ for any $\tau' \in [\tau, \bar{\tau}]$. Now, using Lemma 2, we have $c_2(\alpha) = c_1(\alpha - \tau)$ for any $\alpha \in [\tau - \bar{\tau}, 0]$. Furthermore, if $\bar{c}_2 : \mathfrak{J}(\bar{\mathcal{E}}) \rightarrow S_4^2$ is the maximal integral line of the vector field η^2 with $\bar{c}_2(0) = \bar{\mathcal{E}} = c_1(-\bar{\tau})$, then $\bar{c}_2(\alpha) = c_1(\alpha - \bar{\tau})$ for any $\alpha \in [0, \bar{\tau} - \tau]$. Therefore \bar{c}_2 coincides with

$$\{\alpha \in \mathbf{R} : (\alpha + \tau - \bar{\tau}) \in \mathfrak{J}(\mathcal{E})\} \rightarrow S_4^2, \alpha \mapsto c_2(\alpha + \tau - \bar{\tau}).$$

This yields (i) $\tau \in \mathfrak{J}(\mathcal{E})$ if and only if $\bar{\tau} \in \mathfrak{J}(\bar{\mathcal{E}})$, and (ii) $\bar{c}_2(\bar{\tau}) = c_2(\tau) = \mathcal{E}^2$ for $\tau \in \mathfrak{J}(\mathcal{E})$.

This completes the proof because the case where $\bar{\tau} = \tau$ is trivial, and the one where $\bar{\tau} < \tau$ is similar (it is enough to use the analogue of (1) and Lemma 2 with $\bar{\tau}$ instead of τ and \bar{P}^{21} instead of P^{21}).

We have

$$(\mathcal{E}^1, \mathcal{E}^2), (\tilde{\mathcal{E}}^1, \mathcal{E}^2), (\mathcal{E}^1, \tilde{\mathcal{E}}^2) \in R^{21} \Rightarrow \tilde{\mathcal{E}}^1 = \mathcal{E}^1, \tilde{\mathcal{E}}^2 = \mathcal{E}^2.$$

Therefore R^{21} is the graph of a bijection $f^{21} : \text{Dom } f^{21} \rightarrow \text{Ran } f^{21}$ with

$$(2) \quad (\text{Dom } f^{21}) \cap (\text{Ran } f^{21}) \supset P^{21} \quad \text{and} \quad f^{21}|_{P^{21}} = i_{P^{21}},$$

where $i_{P^{21}} : P^{21} \rightarrow (\text{Ran } f^{21})$ is the inclusion map.

If S_4^2 coincides with S_4^1 , then we can write f^{11} instead of f^{21} and we have

$$(3) \quad f^{11} = i_{S_4^1} \quad \text{— identity map.}$$

In this argument I have used the reflexivity of the divergence relation. Its symmetry allows the exchange of the roles of S_4^1 and S_4^2 . This gives a map f^{12} which satisfies

$$(4) \quad f^{12} = (f^{21})^{-1}.$$

Now, let us use the transitivity, and let us consider S_4^3 in the same equivalence class as S_4^1 and S_4^2 . If $R^{32} [R^{31}]$ is the analogue for $S_4^2 [S_4^1]$ and S_4^3 of the relation R^{21} on S_4^1 to S_4^2 , then

$$(\mathcal{E}^1, \mathcal{E}^2) \in R^{21}, (\mathcal{E}^2, \mathcal{E}^3) \in R^{32} \Rightarrow (\mathcal{E}^1, \mathcal{E}^3) \in R^{31}.$$

But \mathcal{E}^1 can be R^{31} -related to \mathcal{E}^3 also in the case where $(\mathcal{E}^1, \mathcal{E}^2) \notin R^{21}$ and $(\mathcal{E}^2, \mathcal{E}^3) \notin R^{32}$. Thus

$$f^{32} \circ f^{21} : \{\mathcal{E}^1 \in \text{Dom } f^{21} : f^{21}(\mathcal{E}^1) \in \text{Dom } f^{32}\} \rightarrow \text{Ran } f^{32}$$

is generally different from f^{31} . However

$$(5) \quad \mathcal{E}^1 \in \text{Dom } f^{21}, f^{21}(\mathcal{E}^1) \in \text{Dom } f^{32} \Rightarrow \mathcal{E}^1 \in \text{Dom } f^{31}, (f^{32} \circ f^{21})(\mathcal{E}^1) = f^{31}(\mathcal{E}^1).$$

Lastly

PROPOSITION 3. *The map f^{21} is a C^∞ diffeomorphism.*

Proof. If f^{21} is proved to be C^∞ , then f^{12} is C^∞ too, and (4) gives the thesis.

For $k = 1, 2$, assume that (i) $\mathcal{E}^k \in S_4^k$, (ii) $c_{\mathcal{E}^k} : \mathfrak{J}(\mathcal{E}^k) \rightarrow S_4^k$ is the maximal integral curve of the vector field η^k starting from \mathcal{E}^k , (iii) $\tau \in \mathbf{R}$, then we set

$$E_\tau^k = \{\mathcal{E}^k \in S_4^k : \tau \in \mathfrak{J}(\mathcal{E}^k)\} \quad \text{and} \quad \mu_\tau^k : E_\tau^k \rightarrow S_4^k, \mathcal{E}^k \mapsto c_{\mathcal{E}^k}(\tau).$$

(Usually the set $\{\mu_\tau^k : \tau \in \mathbf{R}\}$ is called the flow of the vector field η^k). Since η^k is C^∞ , then E_τ^k is open and μ_τ^k is C^∞ ($\forall \tau \in \mathbf{R}$). By definition, μ_τ^k has a C^∞ inverse which obviously coincides with $\mu_{-\tau}^k$.

Let P^{21} be as in [7] Def. 2, and let us fix (arbitrarily) an event $\mathcal{E}^1 \in \text{Dom } f^{21}$. Moreover let $\tau \in \mathbf{R}$ be such that $\tau \geq 0$ and $c_{\mathcal{E}^1}(-\tau) \in P^{21}$.

Since (i) E_τ^1 and P^{21} are open in S_4^1 , (ii) E_τ^2 is open in S_4^2 , and (iii) the induced topologies of S_4^1 and S_4^2 coincide on P^{21} —see [7] Def. 2—then $P^{21} \cap E_\tau^1 \cap E_\tau^2$ is open, and so is its image $U = \mu_\tau^1(P^{21} \cap E_\tau^1 \cap E_\tau^2)$ under μ_τ^1 , an homeomorphism into.

Now, let us observe that $\mathcal{E}^1 \in U$, $U \subset \text{Dom } f^{21}$, and $f^{21}|_U = (\mu_\tau^2 \circ \mu_{-\tau}^1)|_U$. This implies that f^{21} is C^∞ in the neighbourhood U of the (arbitrary) point \mathcal{E}^1 , because $\mu_{-\tau}^1$ and μ_τ^2 are C^∞ . q.e.d.

* * *

It is interesting to consider the following class \mathfrak{B} of spacetimes.

DEFINITION 4. The spacetime S_4 is said to belong to \mathfrak{B} if $S_4 \in \mathfrak{A}$ —see Def. 3 in ref. [7]—and the reference frame η is a complete vector field—see footnote (5).

The use of \mathfrak{B} instead of \mathfrak{A} in the preceding arguments yields that, in connection with \mathfrak{B} , f^{21} is a C^∞ diffeomorphism from S_4^1 onto S_4^2 , and (5) becomes

$$f^{31} = f^{32} \circ f^{21}.$$

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