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Harmonie reflections

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Geometria differenziale. — *Harmonic reflections.* Nota di LIEVEN VANHECKE e MARIA-ELENA VAZQUEZ-ABAL, presentata (*) dal Socio E. MARTINELLI.

ABSTRACT. — We study local reflections φ_σ with respect to a curve σ in a Riemannian manifold and prove that σ is a geodesic if φ_σ is a harmonic map. Moreover, we prove that the Riemannian manifold has constant curvature if and only if φ_σ is harmonic for all geodesics σ .

KEY WORDS: Harmonic maps; Reflections with respect to a curve; Harmonic reflections.

RIASSUNTO. — *Riflessioni armoniche.* Si studia la riflessione locale φ_σ rispetto ad una curva σ in una varietà riemanniana e si dimostra che σ è una geodetica se φ_σ è un'applicazione armonica. Inoltre si prova che la varietà è a curvatura costante se e solamente se φ_σ è armonica per tutte le geodetiche σ .

1. INTRODUCTION

Local geodesic symmetries on a Riemannian manifold are local diffeomorphisms and the properties of these transformations may be used to characterize several nice classes of Riemannian manifolds. For example, in [3] C. Dodson and the two authors of this paper proved the following: *A Riemannian manifold is locally symmetric if and only if all local geodesic symmetries are harmonic.*

Local geodesic symmetries are local reflections with respect to a point. In [8], [9], [10] this class of transformations has been generalized and the notion of a local reflection with respect to a submanifold, in particular with respect to a curve, has been introduced and studied. These local reflections are local diffeomorphisms. In this paper we consider *harmonic reflections with respect to a curve*. Our aim is to prove

THEOREM 1 — *Let $\sigma: [a, b] \rightarrow M$ be a topologically embedded curve in a Riemannian manifold M . If the local reflection with respect to s is harmonic, then s is a geodesic.*

THEOREM 2 — *Let (M, g) be a connected Riemannian manifold. Then (M, g) is a space of constant curvature if and only if the local reflections with respect to all geodesics are harmonic.*

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2. HARMONIC MAPS

Let (M, g) and (N, h) be two Riemannian manifolds with metrics g and h and let $f: (M, g) \rightarrow (N, h)$ be a smooth map. The pullback f^*h is a semi-definite symmetric covariant tensor of order two, called the *first fundamental form*. Further, the covariant differential $\nabla(df)$ is a symmetric tensor of order two with values in $f^{-1}(TN)$, called the *second fundamental form* of f (see [4], [5]). A map with vanishing second fundamental form is said to be *totally geodesic*.

The trace of $\nabla(df)$ is denoted by $\tau(f)$ and is called the *tension field* of f . A map with vanishing tension field is called a *harmonic map*.

If $U \subset M$ is a domain with coordinates (x^1, \dots, x^m) and $V \subset N$ is a domain with coordinates (y^1, \dots, y^n) such that $f(U) \subset V$, then f can be locally represented by $y^\alpha = f^\alpha(x^1, \dots, x^m)$, $\alpha = 1, \dots, n$. The metric tensor g is represented by $g(x) = g_{ij}(x) dx^i dx^j$, $i, j = 1, \dots, m$, and similarly we have $h(y) = h_{\alpha\beta}(y) dy^\alpha dy^\beta$, $\alpha, \beta = 1, \dots, n$. $df(x)$ is represented by the matrix $\left(\frac{\partial f^\alpha}{\partial x^i}\right)$. In this case we have

$$(f^*h)_{ij} = \frac{\partial f^\alpha}{\partial x^i} \frac{\partial f^\beta}{\partial x^j} h_{\alpha\beta},$$

$$(1) \quad (\nabla(df))_{ij}^\gamma = \frac{\partial^2 f^\gamma}{\partial x^i \partial x^j} - {}^M\Gamma_{ij}^k \frac{\partial f^\gamma}{\partial x^k} + {}^N\Gamma_{\alpha\beta}^\gamma \frac{\partial f^\alpha}{\partial x^i} \frac{\partial f^\beta}{\partial x^j},$$

where ${}^M\Gamma_{ij}^k$ and ${}^N\Gamma_{\alpha\beta}^\gamma$ are the Christoffel symbols for (M, g) and (N, h) respectively.

It follows that f is harmonic if and only if

$$(2) \quad \tau(f)^\gamma = g^{ij} (\nabla(df))_{ij}^\gamma = 0.$$

3. REFLECTIONS WITH RESPECT TO A CURVE

Let $\sigma: [a, b] \rightarrow (M, g)$ be a smooth embedded curve in a Riemannian manifold (M, g) and denote by U a tubular neighbourhood of σ , i.e.

$U = \{p \in M \mid \text{there exists a unique geodesic } \gamma \text{ of } M \text{ through } p \text{ cutting } \sigma \text{ orthogonally}\}$.

Then, for any $p \in U$ we may put

$$p = \exp_{\sigma(t)}(ru), \quad u \in T_{\sigma(t)}^\perp, \quad \|u\| = 1, \quad t \in [a, b],$$

where $r = d(p, \sigma(t))$.

The map $\varphi_\sigma: U \rightarrow U$ defined by

$$\varphi_\sigma: p = \exp_{\sigma(t)}(ru) \mapsto \varphi_\sigma(p) = \exp_{\sigma(t)}(ru)$$

is a local diffeomorphism and is called a *local reflection with respect to* σ .

To treat this map φ_σ analytically we use *Fermi coordinates*. We follow the treat-

ment developed in [7], [9]. For the Riemannian manifold (M, g) let ∇ denote the Levi-Civita connection and R the Riemannian curvature tensor, defined by

$$R_{XY} = \nabla_{[X, Y]} - [\nabla_X, \nabla_Y]$$

for all vector fields X, Y on M .

First, let σ be a unit speed curve, i.e. $\|\dot{\sigma}\| = 1$, and let $\{e_i, i = 1, \dots, n\}$ be an orthonormal basis of $T_{\sigma(a)}M$ such that

$$e_1 = \dot{\sigma}(a).$$

Further, let E_1 be the unit tangent field $\dot{\sigma}$ and E_2, \dots, E_n normal vector fields along σ parallel with respect to the normal connection ∇^\perp of the normal bundle $\dot{\sigma}^\perp$ such that

$$E_i(a) = e_i, \quad i = 2, \dots, n.$$

Then, the Fermi coordinates (x^1, \dots, x^n) with respect to $\sigma(a)$ and (E_1, \dots, E_n) are defined by

$$x^1(\exp_{\sigma(t)} \sum_{j=2}^n t^j E_j) = t - a,$$

$$x^i(\exp_{\sigma(t)} \sum_{j=2}^n t^j E_j) = t^i, \quad \text{where } 2 \leq i \leq n.$$

For a vector $v \in T_{\sigma(t)}^\perp \sigma$ with $\exp_{\sigma(t)} v \in U$ we have

$$v = \sum_{a=2}^n x^a E_a(t) = ru$$

where $\|u\| = 1$ and $r^2 = \sum_{a=2}^n (x^a)^2$.

We put

$$(3) \quad \kappa_u = g(\ddot{\sigma}, u)$$

where $\ddot{\sigma}(t)$ is the (mean) curvature normal of σ at $\sigma(t)$. Note that, when u is parallel with respect to ∇^\perp , we have

$$\nabla_{\dot{\sigma}} u = -\kappa_u \dot{\sigma}.$$

Finally we note that the local reflection φ_σ is now given by

$$\varphi_\sigma : (x^1, x^2, \dots, x^n) \mapsto (x^1, x^2, \dots, x^n).$$

To compute the tension field for φ_σ we need the Christoffel symbols with respect to the system of Fermi coordinates. Therefore we use

$$(4) \quad m_{\Gamma_{ij}^k} = \frac{1}{2} g^{kl} \left\{ \frac{\partial g_{il}}{\partial x^j} + \frac{\partial g_{jl}}{\partial x^i} - \frac{\partial g_{ij}}{\partial x^l} \right\}.$$

So we have to compute g_{ij}, g^{ij} . Therefore we use the method developed in [7] or use Jacobi vector fields as in [6], [9]. The result is given in the following

LEMMA 3. Let $m = \sigma(t)$ and $p = \exp_{\sigma(t)}(su)$, $\|u\| = 1$. With respect to the Fermi coordinates (x^1, \dots, x^n) introduced before, we have

$$(5) \quad g_{11}(p) = \sum_{\ell=0}^4 \alpha_{11}^{\ell}(m) s^{\ell} + O(s^5)$$

where

$$\alpha_{11}^0 = 1,$$

$$\alpha_{11}^1 = -2\kappa_u,$$

$$\alpha_{11}^2 = \kappa_u^2 - R_{1u1u},$$

$$\alpha_{11}^3 = -\frac{1}{3}(\nabla_u R_{1u1u} - 4\kappa_u R_{1u1u}),$$

$$\alpha_{11}^4 = -\frac{1}{12} \left(\nabla R_{uu}^2 R_{1u1u} - 4R_{1u1u}^2 - 4 \sum_{c=2}^n R_{1uc u}^2 - 6\kappa_u \nabla_u R_{1u1u} + 4\kappa_u^2 R_{1u1u} \right);$$

$$(6) \quad g_{1a}(p) = \sum_{\ell=0}^4 \alpha_{1a}^{\ell}(m) s^{\ell} + O(s^5), \quad a = 2, \dots, n,$$

where

$$\alpha_{1a}^0 = \alpha_{1a}^1 = 0,$$

$$\alpha_{1a}^2 = -\frac{2}{3} R_{1uau},$$

$$\alpha_{1a}^3 = -\frac{1}{12} (3 \nabla_u R_{1uau} - 4\kappa_u R_{1uau}),$$

$$\alpha_{1a}^4 = -\frac{1}{30} \left(2\partial_{uu}^2 R_{1uau} - 4R_{1u1u} R_{1uau} - 4 \sum_{c=2}^n R_{1uc u} R_{auc u} - 5\kappa_u \nabla_u R_{1uau} \right);$$

$$(7) \quad g_{ab}(p) = \sum_{\ell=0}^4 \alpha_{ab}^{\ell}(m) s^{\ell} + O(s^5), \quad a, b = 2, \dots, n,$$

where

$$\alpha_{ab}^0 = \delta_{ab},$$

$$\alpha_{ab}^1 = 0,$$

$$\alpha_{ab}^2 = -\frac{1}{3} R_{uaub},$$

$$\alpha_{ab}^3 = -\frac{1}{6} \nabla_u R_{uaub},$$

$$\alpha_{ab}^4 = -\frac{1}{180} \left(9 \nabla_{uu}^2 R_{uaub} - 8R_{1uau} R_{1ubu} - 8 \sum_{c=2}^n R_{uac u}^2 R_{ubuc} \right);$$

$$(8) \quad g^{11}(p) = \sum_{\ell=0}^4 \beta_{11}^{\ell}(m) s^{\ell} + 0(s^5),$$

where

$$\beta_{11}^0 = 1,$$

$$\beta_{11}^1 = 2\kappa_u,$$

$$\beta_{11}^2 = R_{1u1u} + 3\kappa_u^2,$$

$$\beta_{11}^3 = \frac{1}{3}(\nabla_u R_{1u1u} + 8\kappa_u R_{1u1u} + 12\kappa_u^3),$$

$$\beta_{11}^4 = \frac{1}{36} \left(3\nabla_{uu}^2 R_{1u1u} + 24R_{1u1u}^2 + 4 \sum_{c=2}^n R_{1uc}^2 + 30\kappa_u \nabla_u R_{1u1u} + 180\kappa_u^2 R_{1u1u} + 180\kappa_u^4 \right);$$

$$(9) \quad g^{1a}(p) = \sum_{\ell=0}^4 \beta_{1a}^{\ell}(m) s^{\ell} + 0(s^5), \quad a = 2, \dots, n,$$

where

$$\beta_{1a}^0 = \beta_{1b}^1 = 0,$$

$$\beta_{1a}^2 = \frac{2}{3} R_{1uau},$$

$$\beta_{1a}^3 = \frac{1}{4}(\nabla_u R_{1uau} + 4\kappa_u R_{1uau}),$$

$$\beta_{1a}^4 = \frac{1}{45} \left(3\nabla_{uu}^2 R_{1uau} + 24R_{1u1u} R_{1uau} + 4 \sum_{c=2}^n R_{1uc} R_{auc} + 15\kappa_u \nabla_u R_{1uau} + 60\kappa_u^2 R_{1uau} \right);$$

$$(10) \quad g^{ab}(p) = \sum_{\ell=0}^4 \beta_{ab}^{\ell}(m) s^{\ell} + 0(s^5), \quad a, b = 2, \dots, n,$$

where

$$\beta_{ab}^0 = \delta_{ab},$$

$$\beta_{ab}^1 = 0,$$

$$\beta_{ab}^2 = \frac{1}{3} R_{uaub},$$

$$\beta_{ab}^3 = \frac{1}{3} \nabla_u R_{uaub},$$

$$\beta_{ab}^4 = \frac{1}{60} \left(3 \nabla_{uu}^2 R_{uaub} + 24R_{1uau} R_{1ubu} + 4 \sum_{c=2}^n R_{uac} R_{ubuc} \right).$$

4. HARMONIC REFLECTIONS

It follows from (2) and the local expression for the reflection φ_σ (section 3) that φ_σ is harmonic if and only if

$$(11) \quad \tau(\varphi_\sigma)^1(p) = \{g^{11} \nabla(d\varphi_\sigma)_{11}^1 + 2g^{1a} \nabla(d\varphi_\sigma)_{1a}^1 + g^{ab} \nabla(d\varphi_\sigma)_{ab}^1\}(p) = 0,$$

$$(12) \quad \tau(\varphi_\sigma)^c(p) = \{g^{11} \nabla(d\varphi_\sigma)_{11}^c + 2g^{1a} \nabla(d\varphi_\sigma)_{1a}^c + g^{ab} \nabla(d\varphi_\sigma)_{ab}^c\}(p) = 0,$$

for $a, b, c, = 2, \dots, n$, where

$$\nabla(d\varphi_\sigma)_{11}^1(p) = -\Gamma_{11}^1(p) + \Gamma_{11}^1(\varphi_\sigma(p)),$$

$$\nabla(d\varphi_\sigma)_{1a}^1(p) = -\Gamma_{1a}^1(p) \Gamma_{1a}^1(\varphi_\sigma(p)),$$

$$\nabla(d\varphi_\sigma)_{ab}^1(p) = -\Gamma_{ab}^1(p) + \Gamma_{ab}^1(\varphi_\sigma(p)),$$

$$\nabla(d\varphi_\sigma)_{11}^c(p) = \Gamma_{11}^c(p) + \Gamma_{11}^c(\varphi_\sigma(p)),$$

$$\nabla(d\varphi_\sigma)_{1a}^c(p) = \Gamma_{1a}^c(p) \Gamma_{1a}^c(\varphi_\sigma(p)),$$

$$\nabla(d\varphi_\sigma)_{ab}^c(p) = \Gamma_{ab}^c(p) + \Gamma_{ab}^c(\varphi_\sigma(p)).$$

It is worthwhile to note here that $\nabla(d\varphi_\sigma)_{1a}^1$, $\nabla(d\varphi_\sigma)_{11}^c$, $\nabla(d\varphi_\sigma)_{ab}^c$ are even functions and $\nabla(d\varphi_\sigma)_{11}^1$, $\nabla(d\varphi_\sigma)_{ab}^1$, $\nabla(d\varphi_\sigma)_{1a}^c$ are odd functions.

In order to write down the explicit expressions we shall need, we put

$$\frac{\partial g_{ij}}{\partial x^k}(p) = \sum_{\ell=0}^4 \gamma_{ijk}^\ell(m) s^\ell + 0(s^5)$$

and

$$\nabla(d\varphi_\sigma)_{ij}^k(p) = \sum_{\ell=0}^4 A_{ijk}^\ell(m) s^\ell + 0(s^5)$$

for $i, j, k = 1, \dots, n$.

A lengthy computation, which we omit, then leads to

$$(13) \quad \begin{aligned} \tau(\varphi_\sigma)^c(p) = & -\gamma_{11c}^0(m) - \beta_{11}^1(m) \gamma_{11c}^0(m) s + \left(A_{11c}^2 + \sum_{a=2}^n A_{aac}^2 - \gamma_{11c}^0 \beta_{11}^2 \right) (m) s^2 + \\ & + \left(\beta_{11}^3 \gamma_{11c}^0 + 2 \sum_{a=2}^n \beta_{1a}^2 (\gamma_{1ca}^1 - \gamma_{1ac}^1) + \beta_{11}^1 A_{11c}^2 \right) (m) s^3 + 0(s^4) \end{aligned}$$

where

$$(14) \quad \gamma_{11c}^0(m) = -2g(\ddot{\sigma}, E_c)(t).$$

Now we are ready to give the

PROOF OF THEOREM 1 - For a harmonic reflection φ_σ we have $\tau(\varphi_\sigma)^c = 0$, $c = 2, \dots, n$, for all $p \in U$. This implies from (13)

$$\gamma_{11c}^0(m) = \gamma_{11c}^0(\sigma(t)) = 0$$

for all $t \in [a, b]$. Hence, since $\ddot{\sigma}$ is orthogonal to $\dot{\sigma}$, we get from (14) that $\ddot{\sigma} = 0$ which proves the required result.

For the rest of this section we suppose that s is a geodesic. Hence $\kappa_u = 0$ for all u and so

$$(15) \quad v_{11c}^0 = 0, \quad c = 2, \dots, n.$$

Also, from (8), we have

$$(16) \quad \beta_{11}^1 = 0.$$

To prove Theorem 2 we shall use the following

LEMMA 4 - [1], [2]. *Let (M, g) be a connected m -dimensional Riemannian manifold with $m > 2$. Then (M, g) is space of constant curvature if and only if*

$$R(X, Y, X, Z) = 0$$

for any orthogonal triple of vector fields X, Y, Z , on M .

Now we give the

PROOF OF THEOREM 2 - First, let (M, g) be a space of constant curvature. Then, it is proved in [9] it that the local reflections φ_σ with respect to all geodesics σ are isometries. So, since any isometry is a harmonic map, all φ_σ are harmonic.

To prove the converse, we start with $\tau(\varphi_\sigma)^c = 0$, $c = 2, \dots, n$. Using (15) and (16), (13) implies

$$(17) \quad \sum_{a=2}^n \beta_{1a}^2 (\gamma_{1ca}^1 - \gamma_{1ac}^1) = 0.$$

From (9) we have

$$\beta_{1a}^2 = \frac{2}{3} R_{1uau}$$

and from the detailed computations, we obtain

$$\gamma_{1ca}^1 - \gamma_{1ac}^1 = -\frac{2}{3} (R_{1cua} R_{1auc} + 2R_{1uca}).$$

Hence (17) leads to

$$(18) \quad \sum_{a=2}^n R_{1uau} (R_{1cua} - R_{1auc} + 2R_{1uca}) = 0.$$

Now, since u is an arbitrary unit vector normal to the geodesic σ , we may put $u = E_c$. Then (18) implies

$$\sum_{a=2}^n R_{1uau}^2 = 0.$$

Since this must be valid for all geodesics, the required result follows at once from Lemma 4 when $\dim M > 2$.

So we are left with the two-dimensional case. Here, the detailed computations show easily that, with $u = E_2$, we have

$$g_{11}(p) = 1 - s^2 R_{1u1u}(m) - \frac{1}{3} s^3 \nabla_u R_{1u1u}(m) = O(s^4),$$

$$g_{12}(p) = O(s^5),$$

$$g_{22}(p) = 1 + O(s^5),$$

and

$$g^{11}(p) = 1 + s^2 R_{1u1u}(m) + \frac{1}{3} s^3 \nabla_u R_{1u1u}(m) + O(s^4),$$

$$g^{12}(p) = O(s^5),$$

$$g^{22}(p) = 1 + O(s^5).$$

A similar calculation as before now gives

$$\tau(\varphi_\sigma)^2(p) = -s^2(\nabla_u R_{1u1u})(m) + O(s^3)$$

and hence, when φ_σ is harmonic, we must have

$$\nabla_u R_{1u1u} = 0$$

at any point of (M, g) and for all geodesics σ . This implies that the two-dimensional connected manifold has constant curvature.

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