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 Matematica E ApplicazioniRodolfo Salvi

# The equations of viscous incompressible nonhomogeneous fluids in noncylindrical domains: on the existence and regularity 

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Analisi matematica. - The equations of viscous incompressible non-bomogeneous fluids in non-cylindrical domains: On the Existence and Regularity. Nota di Rodolfo Salvi, presentata (*) dal Socio L. Amerio.


#### Abstract

We prove the existence of a weak solution and of a strong solution (locally in time) of the equations which govern the motion of viscous incompressible non-homogeneous fluids. Then we discuss the decay problem.


Key words: Non-homogeneous fluids; Time dependent domains; Weak solutions; Strong solutions.

Riassunto. - Le equazioni dei fluidi viscosi incomprimibili non omogenei in domini non cilindrici: esistenza e regolarità. Si dimostra l'esistenza di una soluzione debole e di una soluzione forte (in piccolo) per le equazioni che governano il moto dei fluidi viscosi incomprimibili con densità non costante. Inoltre si discute il problema dell'andamento asintotico.

## 1. Introduction

We consider the motion of a viscous incompressible non-homogeneous fluid, defined in a domain with moving boundaries. In other words, we have to deal not with a space-time cylinder but with a non-cylindrical domain in $R^{3} \times[0, T]$. To be more precise, we consider a domain $\Omega_{T}=\underset{0 \leqslant t \leqslant T}{ } \Omega(t) \times\{t\}$ where each $\Omega(t)$ is a bounded domain in $R^{3}$, and $T>0$ is a positive number. We will find, in the region $\Omega_{T}$ a solution ( $u, \rho, p$ ) of the system

$$
\begin{equation*}
\rho \partial_{t} u-\mu \Delta u+\rho u \cdot \nabla u+\nabla p-\rho f=0 ; \quad \partial_{t} \rho+u \cdot \nabla_{\rho}=0 ; \quad \nabla \cdot u=0 \quad \text { in } \Omega_{T} \tag{1.1}
\end{equation*}
$$ satisfying the initial-boundary conditions

$$
\begin{cases}u(x, 0)=u_{0} ; & \rho(0)=\rho_{0}  \tag{1.2}\\ u(x, t)=0 & \text { in } \Omega(0), \\ u & \text { on } \Gamma_{T},\end{cases}
$$

where $u=u(t)=\left(u_{1}(x, t), u_{2}(x, t), u_{3}(x, t)\right)$ is the velocity, $\rho=\rho(t)=\rho(x, t)$ the density, $p=p(t)=p(x, t)$ the pressure, $f=f(t)=\left(f_{1}(x, t), f_{2}(x, t), f_{3}(x, t)\right)$ the external force, $\mu$ the viscosity, and $\Gamma_{T}=\underset{0 \leqslant t \leqslant T}{\cup} \Gamma(t) \times\{t\}$ with $\Gamma(t)$ the boundary of $\Omega(t)$.
Problem (1.1), (1.2) was studied in [1,2,5], in cylindrical domains. The paper deals with the existence of weak solutions and of strong (locally in time) solutions of (1.1), (1.2). To prove this, we employ the method developed in [4].

Section 2 contains preliminaries. Section 3 contains the proof of the existence of a weak solution and of a strong (locally in time) solution of (1.1), (1.2), and contains results on the decay problem.
(*) Nella seduta del 21 aprile 1990.

## 2. Preliminaries

All functions in this paper are $R$ - or $R^{3}$-valued. The letter $c$ denotes different constants depending on $\Omega_{T}$ and $\alpha, \beta$ are positive constants. We employ the usual notations of vector analysis; in particular, the $j$-th components of $u \cdot \nabla u$ and $\Delta u$ are

$$
\sum_{i=1}^{3} u_{i} \partial_{x_{i}} u_{j} \quad \text { and } \quad \sum_{i=1}^{3} \partial_{x_{i}} \partial_{x_{i}} u_{j}
$$

respectively. Some additional notation is needed. We let

$$
\begin{gathered}
(u, v)_{\Omega(t)}=\sum_{i=1}^{3} \int_{\Omega(t)} u_{i} v_{i} d x ; \quad|u|_{\Omega(t)}^{2}=(u, u)_{\Omega(t)} ; \quad((u, v))_{\Omega(t)}=\sum_{i=1}^{3} \int_{\Omega(t)} \nabla u_{i} \nabla v_{i} d x ; \\
\|u\|_{\Omega(t)}^{2}=((u, u))_{\Omega(t)} ; \quad|u|_{\Omega_{T}}^{2}=\int_{0}^{T}|u|_{\Omega(t)}^{2} d t ; \quad\|u\|_{\Omega_{T}}^{2}=\int_{0}^{T}\|u\|_{\Omega(t)}^{2} d t ;
\end{gathered}
$$

$$
D(\Omega(t))=\left\{\phi \mid \phi \in\left(C_{0}^{\infty}(\Omega(t))\right)^{3}, \nabla \cdot \phi=0\right\} ; D\left(\Omega_{T}\right)=\left\{\phi \mid \phi \in\left(C^{\infty}\left(\Omega_{T}\right)\right)^{3}, \operatorname{supp} \phi \subset \Omega_{T}, \nabla \cdot \phi=0\right\}
$$

$H(\Omega(t))=$ completion of $D(\Omega(t))$ in the norm $|\phi|_{\Omega(t)} ; V(\Omega(t))=$ completion of $D(\Omega(t))$ in the norm $\|\hat{\phi}\|_{\Omega(t)} ; H\left(\Omega_{T}\right)=$ completion of $D\left(\Omega_{T}\right)$ in the norm $|u|_{\Omega_{T}} ; V\left(\Omega_{T}\right)=$ completion of $D\left(\Omega_{T}\right)$ in the norm $\|u\|_{\Omega_{T}}$. P denotes the projection operator from $L^{2}(\Omega(t))$ onto $H(\Omega(t))$. We assume in the present paper that $\Gamma(t)$ is smooth (at least uniformly of class $C^{3}$ ), and locally represented by the function $x_{3}=\psi\left(x_{1}, x_{2}, t\right)$ (or $x_{1}=\psi\left(x_{2}, x_{3}, t\right)$, or $x_{2}=$ $\left.=\psi\left(x_{3}, x_{1}, t\right)\right)$. Now we are in the position to give the definitions of weak and strong solutions of (1.1), (1.2).
$(u, \rho)$ is a weak solution of (1.1), (1.2) if $u$ and $\rho$ satisfy the following conditions:
(i) $u \in L^{2}(0, T ; V(\Omega(t))) \cap L^{\infty}(0, T ; H(\Omega(t))), \rho \in L^{\infty}\left(\Omega_{T}\right), \quad \alpha \leqslant \rho \leqslant \beta$;
(ii) $\int_{0}^{T}\left\{\left(\rho u, \partial_{t} \phi\right)_{\Omega(t)}+(\rho u \cdot \nabla \phi, u)_{\Omega(t)}-\mu((u, \phi))_{\Omega(t)}+(\rho f, \phi)_{\Omega(t)}\right\} d t=$ $=-\left(\rho_{0} u_{0}, \phi(0)\right)_{\Omega(0)} \quad \forall \phi \in D\left(\Omega_{T}\right)$ with $\phi(T)=0$,
$\partial_{t} \rho+u \cdot \nabla_{\rho}=0 \quad$ in the sense of the distributions;
(iii) $\alpha|u(t)|_{\Omega(t)}^{2}+2 \mu \int_{s}^{t}\|u\|_{\Omega(\sigma)}^{2} d \sigma \leqslant \beta|u(s)|_{\Omega(s)}^{2}+2 \int_{s}^{t}(\rho f, u)_{\Omega(\sigma)} d \sigma$ holds for almost all $s>0$, including $s=0$, and all $t>s$.
$(u, \rho)$ is a strong solution of (1.1), (1.2) if $u$ and $\rho$ satisfy the following conditions:
(i) $u \in L^{2}\left(0, T ; H^{2}(\Omega(t))\right) \cap L^{\infty}(0, T ; V(\Omega(t))), \quad \partial_{t} u \in L^{2}\left(\Omega_{T}\right)$, $p \in L^{\infty}\left(\Omega_{T}\right), \quad \alpha \leqslant \rho \leqslant \beta ;$
(ii) $P\left(\rho \partial_{t} u+\rho u \cdot \nabla u-\mu \Delta u-\rho f\right)=0 \quad$ a.e. in $\Omega_{T}$ $\partial_{t} \rho+u \cdot \nabla_{\rho}=0 \quad$ in the sense of the distributions;
(iii) $|\sqrt{\rho(t)} u(t)|_{\Omega(t)}^{2}+2 \mu \int_{s}^{t}\|u\|_{\Omega(\sigma)}^{2} d \sigma=|\sqrt{\rho(s)} u(s)|_{\Omega(s)}^{2}+2 \int_{s}^{t}(\rho f, u)_{\Omega(\sigma)} d \sigma$
holds for almost all $s>0$, including $s=0$, and all $t>s$.
Our results are now given by the following theorems.
Theorem 1. Let $u_{0} \in H(\Omega(0)), f \in L^{2}\left(\Omega_{T}\right)$, and $\alpha \leqslant \rho_{0} \leqslant \beta$. Then there exists a weak solution of (1.1), (1.2).

Theorem 2. Let $u_{0} \in V(\Omega(t)), f \in L^{2}\left(\Omega_{T}\right)$, and $\alpha \leqslant \rho_{0} \leqslant \beta$. Then there exists a $0<\bar{T} \leqslant T$ such that there exists a strong solution of (1.1), (1.2) in $\Omega_{\bar{T}}$.

Corollary 3. The assumptions of theorem 2 hold. Further, we assume that $\left\|u_{0}\right\|_{\Omega(0)}$ and $|f|_{\Omega_{T}}$ are sufficiently small. Then there exists a strong solution of (1.1), (1.2) for every $T>0$.

Theorem 4. The assumptions of theorem 1 hold, $\Omega(t)$ tends to a bounded domain $\Omega_{0}$ as $t \rightarrow \infty$, and $|P f|_{\Omega(t)} \leqslant c t^{-1 / 2}$. Then there exists a $T_{0}>0$ such that the weak solution of theorem 1 is a strong solution in $\left(T_{0}, \infty\right)$ and $u$ decays like $\|u\|_{\Omega(t)}^{2} \leqslant c t^{-1}$ where $c$ is some positive constant.

Theorem 5. The assumptions of Corollary 3 hold, $\Omega(t)$ tends to a bounded domain $\Omega_{0}$ as $t \rightarrow \infty$, and $|P f|_{\Omega(t)} \leqslant c t^{-1 / 2}$. Then there exists a strong solution of (1.1), (1.2) for every $T>0$ and $u$ decays as $\|u\|_{\Omega_{(t)}}^{2} \leqslant c t^{-1}$ where $c$ is some positive constant.

## 3. Proofs of theorems

First we prove theorem 1 . We consider the following auxiliary problem. Let $\mathfrak{F}=$ $=\left\{\phi \mid \phi \in L^{2}\left(0, T ; H^{2}(\Omega(t))\right), \phi=0\right.$ on $\Gamma_{T}$ with the natural norm $\}$, and $\mathcal{G}=$ $=\left\{\phi \mid \phi \in L^{2}\left(0, T ; H^{2}(\Omega(t))\right), \partial_{t} \phi \in L^{2}\left(0, T ; H^{2}(\Omega(t))\right), \phi=0 \mathrm{on} \Gamma_{T}, \phi(T)=0\right\}$. We consider on $\mathcal{G}$ the norm $\|\phi\|_{\mathcal{G}}=\|\phi\|_{\mathfrak{F}}+\|\phi(0)\|_{\Omega(0)}$.

Find a $\pi_{\varepsilon} \in \mathscr{F}$ such that for all $\phi \in G$,

$$
\begin{align*}
\int_{0}^{T}\left\{\left(\pi_{\varepsilon}, \Delta \partial_{t} \phi\right)_{\Omega(t)}+\varepsilon\left(\Delta \pi_{\varepsilon}, \Delta \phi\right)_{\Omega(t)}-\right. & \left.\left(\tilde{\bar{u}} \cdot \nabla \pi_{\varepsilon}, \Delta \phi\right)_{\Omega(t)}-k\left(\pi_{\varepsilon}, \Delta \phi\right)_{\Omega(t)}\right\} d t=  \tag{3.1}\\
& =\int_{0}^{T} e^{-k t}(w, \Delta \phi)_{\Omega(t)} d t-\left(\left(\rho_{0}-q(0)\right), \Delta \phi\right\}_{\Omega(0)} ;
\end{align*}
$$

here $w=-\partial_{t} q+\varepsilon \Delta q-\widetilde{\bar{u}} \cdot \nabla q ; \bar{u}$ and $q$ are given functions, $k$ is a positive constant, and $\widetilde{\bar{u}}$ is a regularization of $\bar{u}$ through the use of a mollifier depending on a parameter $\lambda$, which is omitted for semplicity.

We denote by $E\left(\pi_{\varepsilon}, \phi\right)$ the left-hand side of (3.1), and by direct computation we have $E(\phi, \phi) \geqslant c_{\varepsilon}\|\phi\|_{G}^{2}$ (for suitable $k$ ); hence by [6, p.208], there exists a $\pi_{\varepsilon} \in \mathfrak{F}$ such that (3.1) holds. Now $-\Delta$ is one to one and onto from $H^{2}(\Omega(t)) \cap H_{0}^{1}(\Omega(t))$ to $L^{2}(\Omega(t))$, then we have that $\pi_{\varepsilon}$ satisfies

$$
\begin{equation*}
\partial_{t} \pi_{\varepsilon}-\varepsilon \Delta \pi_{\varepsilon}+\widetilde{\bar{u}} \cdot \nabla \pi_{\varepsilon}+k \pi_{\varepsilon}=e^{-k t} w \quad \text { a.e. in } \Omega_{T} . \tag{3.2}
\end{equation*}
$$

Now multiplying (3.2) by $\exp k t$ and setting $\rho_{\varepsilon}=\exp k t \pi_{\varepsilon}+q$, we have proved the existence of a solution of the system

$$
\begin{equation*}
\partial_{t} \rho_{\varepsilon}-\varepsilon \Delta \rho_{\varepsilon}+\widetilde{\bar{u}} \cdot \nabla_{\rho_{\varepsilon}}=0 \quad \text { in } \Omega_{T}, \quad \rho_{\varepsilon}=q \text { on } \Gamma_{T} . \tag{3.3}
\end{equation*}
$$

Further, it is a routine matter to prove $\alpha \leqslant \rho_{\varepsilon} \leqslant \beta$.
By using the same method we can prove the existence of a solution $u_{\varepsilon}$ of

$$
\begin{equation*}
P\left(\rho_{\varepsilon} \partial_{t} u_{\varepsilon}-\mu \Delta u_{\varepsilon}+\rho_{\varepsilon} \widetilde{\bar{u}} \cdot \nabla \bar{u}-\rho_{\varepsilon} f-k\left(1-\rho_{\varepsilon}\right)\left(\bar{u}-u_{\varepsilon}\right)\right)=0 . \tag{3.4}
\end{equation*}
$$

Now combining fixed point arguments and a priori estimates proved in [3], passing to limit $\varepsilon \rightarrow 0$, and after $\lambda \rightarrow 0$, we have the existence of a weak solution of (1.1), (1.2).

To prove theorem 2, we consider the approximating system (3.3), (3.4) in which there is not the regularization of $\bar{u}$, in other words we have the terms $\bar{u} \cdot \nabla_{\rho_{\varepsilon}}$ and $\rho_{\varepsilon} \bar{u} \cdot \nabla \bar{u}$ instead of $\widetilde{\bar{u}} \cdot \nabla_{\rho_{\varepsilon}}$ and $\rho_{\varepsilon} \widetilde{\bar{u}} \cdot \nabla \bar{u}$, respectively. Using fixed point arguments, and then passing to limit $\varepsilon \rightarrow 0$, we prove the existence, locally in time, of a strong solution of (1.1), (1.2).

The proof of corollary 3 consists in finding a suitable bound for the data such that fixed point arguments of theorem 2 hold true for every $T>0$. Now combining the estimates of theorem 2 and the results in [4], we can prove theorems $4,5$.

## References

[1] A. V. Kazhikhov, Resolution of boundary value problems for non-bomogeneous fluids. Dokl. Akad. Nauk., 216, 1974, 1008-1010.
[2] O. A. Ladyzenskaya - V. A. Solonnikov, Unique solvability of an initial boundary value problem for the viscous incompressible non-homogeneous fluids. J. Sov. Math., 9, 1978, 697-749.
[3] R. Salvi, On the existence of weak solutions of a non-linear mixed problem for non-homogeneous fluids in a time dependent domain. C.M.U.C., 26, 1985, 185-199.
[4] R. Salvi, On the Navier-Stokes equations in non-cylindrical domains: On the Existence and Regularity. Math. Z., 199, 1988, 153-170.
[5] R. Salvi, The equations of viscous incompressible non-homogeneous fluid: On the Existence and Regularity. Submitted to J. Australian Math. Soc.
[6] F. Treves, Basic linear partial differential equations. Academic Press, New York 1975.

