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## TANGENTIAL CAUCHY-RIEMANN EQUATIONS ON QUADRATIC CR MANIFOLDS

ABSTRACT. — We study the tangential Cauchy-Riemann equations  $\bar{\partial}_b u = \omega$  for  $(0, q)$ -forms on quadratic CR manifolds. We discuss solvability for data  $\omega$  in the Schwartz class and describe the range of the tangential Cauchy-Riemann operator in terms of the signatures of the scalar components of the Levi form.

KEY WORDS: Tangential Cauchy-Riemann complex; Kohn Laplacian; CR manifolds; Global solvability; Hypoellipticity.

### 1. INTRODUCTION

Let  $V$  be an  $n$ -dimensional complex vector space,  $W$  an  $m$ -dimensional real vector space,  $W^{\mathbb{C}}$  the complexification of  $W$ , and

$$\Phi : V \times V \longrightarrow W^{\mathbb{C}}$$

a Hermitian map (i.e.  $\Phi(z, z') = \overline{\Phi(z', z)}$  for every  $z, z' \in V$ , where complex conjugation in  $W^{\mathbb{C}}$  is referred to the real form  $W$ ).

We consider the associated *quadratic manifold*

$$(1) \quad S = \{(z, t + iu) \in V \times W^{\mathbb{C}} : u = \Phi(z, z)\}$$

in  $n + m$  complex dimensions. Then  $S$  is a CR manifold of CR-dimension  $n$  and real codimension  $m$ .

We consider the  $\bar{\partial}_b$ -complex on  $S$ , mapping  $(0, q)$ -forms on  $S$  into  $(0, q + 1)$ -forms, for  $0 \leq q \leq n$ .

We shall consistently use the parameters  $(z, t) \in V \times W$  to denote the element  $(z, t + i\Phi(z, z)) \in S$ . A natural Lie group structure can be introduced on  $V \times W$  (as described in Section 1); this group will be denoted by  $G_{\Phi}$ .

The fiber of the vector bundle  $\Lambda^{0,q}(T^*S)$  over each point of  $S$  can be identified in the trivial way with the exterior product  $\Lambda_q = \Lambda^{0,q}(V^*)$ . Through the identification of  $S$  with  $V \times W = G_{\Phi}$ , we then regard  $(0, q)$ -forms on  $S$  as vector valued functions on  $G_{\Phi}$  with values in  $\Lambda_q$ .

Depending on the integrability or regularity conditions imposed on the forms under consideration, we shall denote the different spaces of  $(0, q)$ -forms as  $L^2(G_{\Phi}) \otimes \Lambda_q$ ,  $S(G_{\Phi}) \otimes \Lambda_q$ ,  $S'(G_{\Phi}) \otimes \Lambda_q$ , etc.

We shall also need other linear bundles over  $G_{\Phi}$ , with fibers  $\text{End}(\Lambda_q)$ ,  $\text{Hom}(\Lambda_q, \Lambda_{q+1})$ , etc. The corresponding spaces of sections will be denoted in a similar way.

We address the following problem: determine under which assumptions on the Hermitean form  $\Phi$ ,  $q \geq 1$  and the  $\bar{\partial}_b$ -closed form  $\omega \in \mathcal{S}(G_\Phi) \otimes \Lambda_q$ , the equation  $\bar{\partial}_b u = \omega$  has a solution  $u \in \mathcal{S}'(G_\Phi) \otimes \Lambda_q$ .

Our approach is based on the results in [7] concerning the Kohn Laplacian

$$\square_b^{(q)} = \bar{\partial}_b \bar{\partial}_b^* + \bar{\partial}_b^* \bar{\partial}_b$$

acting on  $(0, q)$ -forms on  $S$ . The operator  $\bar{\partial}_b^*$  is the adjoint of the  $L^2$ -closure of  $\bar{\partial}_b$  w.r. to the Haar measure  $dz dt$  on  $G_\Phi$  and to a fixed inner product on  $V$ .

The answer depends on the signatures of the scalar-valued forms

$$\Phi^\lambda(z, z') = \lambda(\Phi(z, z')) ,$$

depending on  $\lambda \in W^*$ .

In general,  $W^*$  decomposes as the union of an open regular set, consisting of those  $\lambda$  for which  $\Phi^\lambda$  is non-degenerate, and its complement, the singular set. We further decompose the regular set as the union of the open sets  $\Omega_q$ , defined by the condition that  $\Phi^\lambda$  has  $q$  positive and  $n - q$  negative eigenvalues. Some  $\Omega_q$  may be empty, and it may well happen that all of them are empty, i.e. that  $\Phi^\lambda$  is degenerate for every  $\lambda$ .

In general, we denote by  $\Omega \subset W^*$  the set of those  $\lambda$  for which  $\Phi^\lambda$  has maximum rank. If there are non-degenerate  $\Phi^\lambda$ , then  $\Omega$  is the regular set.

In [7] we proved the following theorem.

**THEOREM 1.1.** *The following are equivalent :*

- (i)  $\Omega_q$  is non-empty;
- (ii)  $\square_b^{(q)}$  is locally solvable, i.e. given any smooth  $(0, q)$ -form  $\omega$ , the equation  $\square_b^{(q)} u = \omega$  has a solution in a fixed neighborhood of the origin;
- (iii)  $\square_b^{(q)}$  has a tempered fundamental solution, i.e.  $K_q \in \mathcal{S}'(G_\Phi) \otimes \text{End}(\Lambda_q)$  such that  $\square_b^{(q)}(\omega * K_q) = \omega$  for every  $\omega \in \mathcal{S}(G_\Phi) \otimes \Lambda_q$ ;
- (iv) the  $L^2$ -null-space of  $\square_b^{(q)}$  is trivial.

If  $\Omega_q$  is non-empty, then  $\square_b^{(q)}$  has a relative fundamental solution  $K_{q,rel} \in \mathcal{S}'(G_\Phi) \otimes \text{End}(\Lambda_q)$ , i.e. such that  $\square_b^{(q)}(\omega * K_{q,rel}) = (I - C_q)\omega$  for every  $\omega \in \mathcal{S}(G_\Phi) \otimes \Lambda_q$ , where  $C_q$  is the orthogonal projection of  $L^2(G_\Phi) \otimes \Lambda_q$  onto the null-space of  $\square_b^{(q)}$ .

The notation we have used is such that, if  $f$  and  $g$  are functions on  $G_\Phi$  with values in  $\Lambda_q$  and in  $\text{End}(\Lambda_q)$  respectively, then

$$f * g(z, t) = \int_{G_\Phi} g((w, u)^{-1}(z, t))f(w, u) dw du$$

takes values in  $\Lambda_q$ .

We derive from this the following result.

**THEOREM 1.2.** *Let  $q \geq 1$ . The equation  $\bar{\partial}_b u = \omega$  has a solution  $u \in \mathcal{S}'(G_\Phi) \otimes \Lambda_{q-1}$  for a given  $\bar{\partial}_b$ -closed form  $\omega \in \mathcal{S}(G_\Phi) \otimes \Lambda_q$  if and only if  $\mathcal{C}_q \omega = 0$ . In particular, the equation has a solution  $u \in \mathcal{S}'(G_\Phi) \otimes \Lambda_{q-1}$  for every  $\omega \in \mathcal{S}(G_\Phi) \otimes \Lambda_q$  such that  $\bar{\partial}_b \omega = 0$  if and only if  $\Omega_q$  is empty.*

The proof will require a precise description of the  $L^2$ -null-space of  $\square_b^{(q)}$ . This will be done in Section 3.

The question of solvability for the Cauchy-Riemann complex has drawn a great deal of interest since Lewy’s celebrated example of a non-solvable differential operator [6]. Such question is also of interest for extension phenomena, such as Bochner’s theorem, see [1-5] for historical background and references.

The relations between solvability of the  $\bar{\partial}_b$ -complex and signatures of the scalar components of the Levi form for general CR manifolds have been investigated by several authors [9-11].

In the case of a quadratic CR manifold  $S$ , Rossi and Vergne [8] showed that if  $\Phi$  is non-degenerate, *i.e.* there exists  $\lambda \in W^*$  such that  $\Phi^\lambda$  is non-degenerate, then condition (i) in Theorem 1.1 is necessary and sufficient for the solvability of the  $\bar{\partial}_b$ -equation. The degenerate case is not included in their analysis. On the other hand, it is as an immediate consequence of our Theorem 1.2 that if  $\Phi$  is degenerate, then  $\Omega_q$  is empty and the  $\bar{\partial}_b$ -equation  $\bar{\partial}_b u = \omega$  is solvable for all  $\bar{\partial}_b$ -closed,  $(0, q)$ -forms  $\omega$  in the Schwartz class, for all  $q \geq 1$ .

2. THE LIE GROUP ASSOCIATED TO  $S$  AND ITS REPRESENTATIONS

We define the following product between two elements  $(z, t), (z', t') \in V \times W$ :

$$(z, t)(z', t') = (z + z', t + t' + 2\text{Im } \Phi(z, z')) ,$$

which induces a step-two nilpotent Lie group structure on  $V \times W$ . We call  $G_\Phi$  this group and  $\mathfrak{g}_\Phi$  its Lie algebra.

For  $v \in V$ , let  $\partial_v f$  denote the directional derivative of a function  $f$  in the direction  $v$ . The left-invariant vector field  $X_v$  on  $G_\Phi$  that coincides with  $\partial_v$  at the origin is given by

$$X_v f(z, t) = \partial_v f(z, t) + 2\text{Im } \Phi(z, v) \cdot \nabla_t f(z, t) .$$

If  $J$  denotes the complex structure on  $V$ , we define  $Z_v, \bar{Z}_v \in \mathfrak{g}_\Phi^{\mathbb{C}}$  as

$$\begin{aligned} Z_v &= \frac{1}{2}(X_v - iX_{Jv}) = \frac{1}{2}(\partial_v - i\partial_{Jv}) + i\overline{\Phi(z, v)} \cdot \nabla_t , \\ \bar{Z}_v &= \frac{1}{2}(X_v + iX_{Jv}) = \frac{1}{2}(\partial_v + i\partial_{Jv}) - i\Phi(z, v) \cdot \nabla_t . \end{aligned}$$

The relevance of the group  $G_\Phi$  in our context is justified by the fact that the operators  $\bar{Z}_v$  coincide with the tangential Cauchy-Riemann operators on  $S$ .

The following commutation rules hold:

$$(2) \quad \begin{aligned} [Z_v, Z_{v'}] &= [\bar{Z}_v, \bar{Z}_{v'}] = 0, \\ [Z_v, \bar{Z}_{v'}] &= -2i\Phi(v, v') \cdot \nabla_{t'}. \end{aligned}$$

Hence  $G_\Phi$  is step-two nilpotent.

It can be shown that the Lie algebras that arise in this way can be characterized as follows.

PROPOSITION 2.1. *A real Lie algebra  $\mathfrak{g}$  is isomorphic to the Lie algebra  $\mathfrak{g}_\Phi$  associated to a quadratic CR manifold if and only if  $\mathfrak{g} = \mathfrak{v} \oplus \mathfrak{w}$ , with  $[\mathfrak{v}, \mathfrak{v}] \subseteq \mathfrak{w}$ ,  $[\mathfrak{g}, \mathfrak{w}] = 0$  and there is a complex structure  $J$  on  $\mathfrak{v}$  such that  $[Jv, Jv'] = [v, v']$  for every  $v, v' \in \mathfrak{v}$ .*

We summarize the description of the irreducible unitary representations of  $G_\Phi$ .

For  $\lambda \in W^*$ , let  $\Phi^\lambda$  be the scalar-valued form  $\Phi^\lambda(v, v') = \lambda(\Phi(v, v'))$ . Denote by  $V_0^\lambda$  the radical of  $\Phi^\lambda$ , i.e. the subspace of the  $v$  such that  $\Phi^\lambda(v, v') = 0$  for every  $v' \in V$ .

Let  $V_1^\lambda$  be the orthogonal complement of  $V_0^\lambda$  in  $V$ , w.r. to the fixed inner product. Let also  $V_r^\lambda$  be a real form of  $V_1^\lambda$  on which  $\Phi^\lambda$  is real. For  $z \in V$ , we set  $z = z' + z''$ , with  $z' \in V_1^\lambda$ ,  $z'' \in V_0^\lambda$ , and  $z' = x' + iy'$  with  $x', y' \in V_r^\lambda$ .

LEMMA 2.2. *Let  $\lambda \in W^*$ , and let  $\eta$  be a linear functional on  $V_0^\lambda$ . Define the representation  $\pi_{\lambda, \eta}$  of  $G_\Phi$  on  $L^2(V_r^\lambda)$  as*

$$(3) \quad (\pi_{\lambda, \eta}(z, t)\phi)(\xi) = e^{i(\lambda(t) + 2\text{Re } \eta(z''))} e^{-2i\Phi^\lambda(y', \xi + x')} \phi(\xi + 2x').$$

Then  $\pi_{\lambda, \eta}$  is an irreducible unitary representation of  $G_\Phi$ . Conversely, any irreducible unitary representation of  $G_\Phi$  is equivalent to one and only one  $\pi_{\lambda, \eta}$ .

It is convenient to diagonalize  $\Phi^\lambda$  with respect to an orthonormal basis  $\{v_1^\lambda, \dots, v_n^\lambda\}$  of  $V$ , in such a way that  $v_j^\lambda \in V_r^\lambda$  for  $j \leq \nu(\lambda)$  and  $v_j^\lambda \in V_0^\lambda$  for  $j > \nu(\lambda)$ , where  $0 \leq \nu(\lambda) = \text{rank } \Phi^\lambda = \dim V_1^\lambda \leq n$ . We set

$$(4) \quad \mu_j = \mu_j(\lambda) = \Phi^\lambda(v_j^\lambda, v_j^\lambda).$$

Calling

$$Z_j^\lambda = \frac{1}{2}(X_{v_j^\lambda} - iX_{Jv_j^\lambda}), \quad \bar{Z}_j^\lambda = \frac{1}{2}(X_{v_j^\lambda} + iX_{Jv_j^\lambda}),$$

a standard computation gives that

$$(5) \quad \begin{aligned} d\pi_{\lambda, \eta}(Z_k^\lambda) &= \begin{cases} \partial_{\xi_j} - \mu_k \xi_k & \text{if } k \leq \nu \\ i\bar{\eta}_{k-\nu(\lambda)} & \text{if } k > \nu \end{cases} \\ d\pi_{\lambda, \eta}(\bar{Z}_k^\lambda) &= \begin{cases} \partial_{\xi_j} + \mu_k \xi_k & \text{if } k \leq \nu \\ i\eta_{k-\nu(\lambda)} & \text{if } k > \nu, \end{cases} \end{aligned}$$

with  $\eta_{k-\nu(\lambda)} = \eta(v_k^\lambda) \in \mathbb{C}$ .

For a function  $f$  on  $G_\Phi$ , we define

$$(6) \quad \pi_{\lambda, \eta}(f) = \int f(z, t) \pi_{\lambda, \eta}(z, t)^{-1} dz dt .$$

This definition has the effect that  $\pi_{\lambda, \eta}(f * g) = \pi_{\lambda, \eta}(g) \pi_{\lambda, \eta}(f)$ , and that

$$\pi_{\lambda, \eta}(\mathcal{L}f) = d\pi_{\lambda, \eta}(\mathcal{L})\pi_{\lambda, \eta}(f) ,$$

for any left-invariant differential operator  $\mathcal{L}$ .

We denote by  $h_j$  the  $j$ -th Hermite function on the real line:

$$(7) \quad h_j(t) = (2^j \sqrt{\pi} j!)^{-1/2} (-1)^j e^{t^2/2} \frac{d^j}{dt^j} e^{-t^2} ,$$

and, for a given a multi-index  $m \in \mathbb{N}^{\nu(\lambda)}$ , we set

$$(8) \quad h_m^\lambda(\xi) = \prod_{j=1}^{\nu(\lambda)} |\mu_j|^{1/4} h_{m_j}(|\mu_j|^{1/2} \xi_j) .$$

By (7) and (8) above we have that  $(d_t - t)h_j = \sqrt{2(j+1)}h_{j+1}$  and  $(d_t + t)h_j = -\sqrt{2j}h_{j-1}$ . From these it follows that

$$(\partial_{\xi_k} + \mu_k \xi_k) h_m^\lambda = (-\text{sgn } \mu_k) ((2m_k + 1)|\mu_k| - \mu_k)^{1/2} h_{m - (\text{sgn } \mu_k) e_k}^\lambda ,$$

where  $e_k = (0, \dots, 1, \dots, 0)$  denotes the  $k$ -th element of the standard basis. Hence, for any unitary irreducible representation  $\pi_{\lambda, \eta}$  of  $G_\Phi$ ,

$$d\pi_{\lambda, \eta}(\bar{Z}_k^\lambda) h_m^\lambda = \begin{cases} (-\text{sgn } \mu_k) ((2m_k + 1)|\mu_k| - \mu_k)^{1/2} h_{m - (\text{sgn } \mu_k) e_k}^\lambda & \text{if } k \leq \nu(\lambda) \\ i\eta_{k - \nu(\lambda)} h_m^\lambda & \text{if } k > \nu(\lambda) . \end{cases}$$

Analogously, one gets

$$d\pi_{\lambda, \eta}(Z_k^\lambda) h_m^\lambda = \begin{cases} \text{sgn } \mu_k ((2m_k + 1)|\mu_k| + \mu_k)^{1/2} h_{m + (\text{sgn } \mu_k) e_k}^\lambda & \text{if } k \leq \nu(\lambda) \\ i\bar{\eta}_{k - \nu(\lambda)} h_m^\lambda & \text{if } k > \nu(\lambda) . \end{cases}$$

The matrix coefficient  $\langle \pi_{\lambda, \eta}(f) h_m^\lambda, h_{m'}^\lambda \rangle$  will be denoted as  $\widehat{f}(\lambda, \eta; m, m')$ . For a  $(0, q)$ -form  $\omega$ , the notation  $\widehat{\omega}(\lambda, \eta; m, m') \in \Lambda_q$  will be used with the same meaning. If  $\Phi^\lambda$  is non-degenerate, so that  $V_0^\lambda = \{0\}$ , we drop the parameter  $\eta$ .

Define

$$D(\lambda) = \prod_{j=1}^{\nu(\lambda)} |\mu_j| .$$

LEMMA 2.3. *The function  $D(\lambda)$  is smooth on  $\Omega$ , the subset of  $W^*$  where  $\nu(\lambda)$  is maximum. The Plancherel formula for  $G_\Phi$  is*

$$(9) \quad \|f\|_2^2 = \int_\Omega \int_{(V_0^\lambda)^*} \|\pi_{\lambda, \eta}(f)\|_{HS}^2 d\eta D(\lambda) d\lambda ,$$

where  $d\lambda$  is an appropriately normalized Lebesgue measure on  $W^*$  and  $d\eta$  is the volume element on  $(V_0^\lambda)^*$  induced by the inner product on  $V$ .

Observe that the domain of integration in (9) has a natural differentiable structure, as it can be identified with  $\Omega \times \mathbb{C}^{n-\nu}$ , with  $\nu = \max \nu(\lambda)$ .

3. THE  $\bar{\partial}_b$ -COMPLEX ON  $G_\Phi$  AND THE NULL SPACE OF  $\square_b^{(q)}$

Let  $\{v_1, \dots, v_n\}$  be any orthonormal basis of  $V$  with respect to the given inner product. Let  $(z_1, \dots, z_n)$  denote the coordinates on  $V$  with respect to this basis. We denote by  $d\bar{z}^I$  the  $(0, q)$ -form  $d\bar{z}_{i_1} \wedge \dots \wedge d\bar{z}_{i_q}$ , where  $I = (i_1, \dots, i_q)$  is a strictly increasing multi-index. Given a  $(0, q)$ -form  $\phi = \sum_{|I|=q} \phi_I d\bar{z}^I$  with smooth coefficients, we have

$$(10) \quad \bar{\partial}_b \phi = \sum_{|I|=q} \sum_{k=1}^n \bar{Z}_k(\phi_I) d\bar{z}_k \wedge d\bar{z}^I = \sum_{|I|=q+1} \sum_{k, |I|=q} \epsilon_{kl}^I \bar{Z}_k(\phi_I) d\bar{z}^I,$$

where

$$Z_j = \frac{1}{2}(X_{v_j} - iX_{Jv_j}), \quad \bar{Z}_j = \frac{1}{2}(X_{v_j} + iX_{Jv_j}), \quad j = 1, \dots, n,$$

and  $\epsilon_{kl}^J = 0$  if  $J \neq \{k\} \cup I$  as sets, and it equals the parity of the permutation that rearranges  $(k, i_1, \dots, i_q)$  in increasing order if  $J = \{k\} \cup I$ .

Then  $\bar{\partial}_b^*$  can be easily computed to yield that

$$(11) \quad \bar{\partial}_b^* \left( \sum_{|I|=q} \phi_I d\bar{z}^I \right) = \sum_{|I|=q-1} \left( - \sum_{k, |I|=q} \epsilon_{kj}^I Z_k \phi_I \right) d\bar{z}^I.$$

The Kohn Laplacian is defined as  $\square_b^{(q)} = \bar{\partial}_b \bar{\partial}_b^* + \bar{\partial}_b^* \bar{\partial}_b$ . This is explicitly computed in [7, Proposition 2.1].

We assume now that  $\Omega_q$  is non-empty. Then also  $\Omega_{n-q} = -\Omega_q$  is non-empty and it contains a Zariski-open subset  $\Omega'_{n-q}$  where the number of distinct eigenvalues of  $\Phi^\lambda$  is maximum. It is shown in [7] that locally on  $\Omega'_{n-q}$  the eigenvalues  $\mu_j(\lambda)$  and the basis elements  $v_j^\lambda$  in (4) are well-defined real-analytic functions of  $\lambda$ .

We can therefore find a locally finite open covering  $\{U_j\}$  of  $\Omega'_{n-q}$  such that for each  $j$  there is an orthonormal coordinate system  $(z_1^\lambda, \dots, z_n^\lambda)$  on  $V$  that varies smoothly with  $\lambda \in U_j$  and diagonalizes  $\Phi^\lambda$  as  $\Phi^\lambda(z, z) = \sum_{k=1}^n \mu_k |z_k^\lambda|^2$ .

For a multi-index  $L$  of length  $q$  with entries  $\ell_1 < \ell_2 < \dots < \ell_q$ , we denote by  $\omega_L^\lambda$  the form  $d\bar{z}_{\ell_1}^\lambda \wedge \dots \wedge d\bar{z}_{\ell_q}^\lambda$ .

Let  $\bar{L} = \bar{L}_j$  the multi-index of length  $q$  containing those  $k$  for which  $\mu_k < 0$ . Let also  $\{\rho_j\}$  be a smooth partition of unity on  $\Omega'_{n-q}$  subordinated to the given covering.

The following result is proven in [7, Lemma 5.1].



LEMMA 3.1. *Let  $\omega \in L^2(G_{\mathbb{F}}) \otimes \Lambda_q$ . The following are equivalent:*

- (i)  $\omega$  is in the null space of  $\square_b^{(q)}$ ;
- (ii)  $\pi_\lambda(\omega) = 0$  a.e. outside of  $\Omega_{n-q}$  and, a.e. on each  $U_j$ ,  $\pi_\lambda(\omega) = T^\lambda \otimes \omega_T^\lambda$ , where  $T^\lambda$  is a Hilbert-Schmidt operator on  $L^2(V_r^\lambda)$ , with range in the linear span of  $b_0^\lambda$ .

COROLLARY 3.2. *The subspace  $(\mathcal{S}(G_{\mathbb{F}}) \otimes \Lambda_q) \cap \ker \square_b^{(q)}$  is dense in  $(L^2(G_{\mathbb{F}}) \otimes \Lambda_q) \cap \ker \square_b^{(q)}$  in the  $L^2$ -topology. Moreover, if  $\omega \in (L^2(G_{\mathbb{F}}) \otimes \Lambda_q) \cap \ker \square_b^{(q)}$  then  $\bar{\partial}_b \omega = \bar{\partial}_b^* \omega = 0$  in the sense of distributions.*

PROOF. By Lemma 5.1 in [7] it follows that  $\omega \in L^2(G_{\mathbb{F}}) \otimes \Lambda_q$  lies in  $\ker \square_b^{(q)}$  if and only if  $\widehat{\omega}(\lambda; k, \ell) \neq 0$  implies that  $\lambda \in \Omega_{n-q}$ ,  $\ell = 0$  and  $\widehat{\omega}(\lambda; k, 0) = c(\lambda, k)\omega_T^\lambda$  is such that

$$\|\omega\|_{L^2}^2 = \int_{\Omega_{n-q}} \sum_k |c(\lambda, k)|^2 D(\lambda) d\lambda.$$

Therefore, for any  $\varepsilon > 0$  it is possible to find a positive integer  $k_0$  and Schwartz functions  $\psi_k$  with support in  $\Omega_{n-q}$ , identically zero for  $k > k_0$ , and such that

$$\int_{\Omega_{n-q}} \sum_k |c(\lambda, k) - \psi_k(\lambda)|^2 D(\lambda) d\lambda < \varepsilon.$$

By Lemmas 3.1 and 5.2 in [7] there exists  $\psi \in \mathcal{S}(G_{\mathbb{F}}) \otimes \Lambda_q$  such that  $\widehat{\psi}(\lambda; k, \ell) = \delta_{0\ell} \psi_k(\lambda)\omega_T^\lambda$ . Hence  $\psi \in \ker \square_b^{(q)}$  and  $\|\omega - \psi\| < \varepsilon$ .

The second assertion is clear for Schwartz forms and follows from the density above for an  $L^2$ -form.  $\square$

#### 4. PROOF OF THEOREM 1.2

The proof is based on the following lemma.

LEMMA 4.1. *There is a family  $\{K_q\}_{0 \leq q \leq n}$ , with  $K_q \in \mathcal{S}'(G_{\mathbb{F}}) \otimes \text{End}(\Lambda_q)$ , satisfying the following properties*

- (i)  $K_q$  is a fundamental (resp. a relative fundamental) solution of  $\square_b^{(q)}$  if  $\Omega_q$  is empty (resp. non-empty);
- (ii) the following identity holds

$$(12) \quad \bar{\partial}_b(\omega * K_q) = (\bar{\partial}_b \omega) * K_{q+1},$$

for all  $\omega \in \mathcal{S}(G_{\mathbb{F}}) \otimes \Lambda_q$ .

Assuming the validity of the lemma, we prove Theorem 1.2.

PROOF OF THEOREM 1.2. Given  $\omega$  as in the statement, it suffices to define  $u = \bar{\partial}_b^*(\omega * K_q)$ . Since  $\bar{\partial}_b \omega = 0$ , we have

$$\bar{\partial}_b^* \bar{\partial}_b(\omega * K_q) = \bar{\partial}_b^*(\bar{\partial}_b \omega * K_{q+1}) = 0,$$

by (12). Then

$$\bar{\partial}_b u = \bar{\partial}_b \bar{\partial}_b^* (\omega * K_q) = \square_b^{(q)} (\omega * K_q) = (I - \mathcal{C}_q) \omega.$$

On the other hand, notice that if  $\omega \in \mathcal{S}(G_\Phi) \otimes \Lambda_q$  is such that  $\mathcal{C}_q \omega \neq 0$ , then the equation  $\bar{\partial}_b u = \omega$  cannot be solved. Indeed, using Corollary 3.2, let  $\omega_k \in (\mathcal{S}(G_\Phi) \otimes \Lambda_q) \cap \ker \square_b^{(q)}$  be a sequence converging to  $\mathcal{C}_q \omega$  in  $L^2(G_\Phi) \otimes \Lambda_q$ . Then,

$$\|\mathcal{C}_q \omega\|_{L^2}^2 = \langle \omega, \mathcal{C}_q \omega \rangle = \lim_{k \rightarrow +\infty} \langle \bar{\partial}_b u, \omega_k \rangle = \lim_{k \rightarrow +\infty} \langle u, \bar{\partial}_b^* \omega_k \rangle = 0,$$

a contradiction.  $\square$

PROOF OF LEMMA 4.1. Given  $\lambda \in \Omega$  and (in case  $\nu = \text{rank } \Phi^\lambda < n$ )  $\eta \in (V_0^\lambda)^*$ , define the operator  $A_q^{\lambda, \eta}$  on  $L^2(V_r^\lambda) \otimes \Lambda_q$  as

$$A_q^{\lambda, \eta} (b_m^\lambda \otimes \omega_L^\lambda) = \begin{cases} 0 & \text{if } \nu = n, \lambda \in \Omega_q \\ & m = 0 \text{ and } L = \bar{L} \\ \frac{1}{\alpha_L^\lambda + |\eta|^2 + \sum_{j=1}^n (2m_j + 1) |\mu_j|} (b_m^\lambda \otimes \omega_L^\lambda) & \text{otherwise.} \end{cases}$$

When  $\nu = n$ , we simply drop  $\eta$  from this formula altogether. Due to its diagonal form, it is easy to check that  $A_q^{\lambda, \eta}$  is a bounded operator.

The (relative) fundamental solutions  $K_q$  constructed in [7] are such that for any pair of Schwartz  $(0, q)$ -forms  $\omega, \sigma$  on  $G_\Phi$ ,

$$(13) \quad \langle \omega * K_q, \sigma \rangle = - \int_{\Omega} \int_{(V_0^\lambda)^*} \langle A_q^{\lambda, \eta} \pi_{\lambda, \eta}(\omega), \pi_{\lambda, \eta}(\sigma) \rangle d\eta D(\lambda) d\lambda.$$

The inner product  $\langle \cdot, \cdot \rangle$  is the ordinary Hilbert-Schmidt inner product for operators on  $L^2(V_r^\lambda) \otimes \Lambda_q$ . If  $\nu = n - 1$ , the integral in  $d\eta$  in (13) may not be absolutely convergent for certain values of  $\lambda$ , and it must be taken in a principal value sense.

We shall show below that

$$(14) \quad d\pi_{\lambda, \eta}(\bar{\partial}_b) \circ A_q^{\lambda, \eta} = A_{q+1}^{\lambda, \eta} \circ d\pi_{\lambda, \eta}(\bar{\partial}_b).$$

This implies that, if  $\omega \in \mathcal{S}(G_\Phi) \otimes \Lambda_q$  and  $\sigma \in \mathcal{S}(G_\Phi) \otimes \Lambda_{q+1}$ , then

$$(15) \quad \langle A_q^{\lambda, \eta} \pi_{\lambda, \eta}(\omega), \pi_{\lambda, \eta}(\bar{\partial}_b^* \sigma) \rangle = \langle A_{q+1}^{\lambda, \eta} \pi_{\lambda, \eta}(\bar{\partial}_b \omega), \pi_{\lambda, \eta}(\sigma) \rangle.$$

Hence

$$(16) \quad \langle \omega * K_q, \bar{\partial}_b^* \sigma \rangle = \langle (\bar{\partial}_b \omega) * K_q, \sigma \rangle,$$

which implies (12). When  $\nu = n - 1$ , the derivation of (16) from (15) requires some more care, but we leave the details to the interested reader.

The proof of (14) is very easy when  $\nu < n$  and  $\eta \neq 0$ , or when  $\nu = n$  and  $\lambda \notin \Omega_q \cup \Omega_{q+1}$ . In both cases, in fact,  $A_q^{\lambda, \eta} = -d\pi_{\lambda, \eta}(\square_b^{(q)})^{-1}$  and  $A_{q+1}^{\lambda, \eta} = -d\pi_{\lambda, \eta}(\square_b^{(q+1)})^{-1}$ . It is sufficient to apply  $d\pi_{\lambda, \eta}$  to both sides of the identity  $\bar{\partial}_b \square_b^{(q)} = \square_b^{(q+1)} \bar{\partial}_b$ .

Assume now that  $\nu = n$  and  $\lambda \in \Omega_q$ . We recall that

$$d\pi_\lambda(\bar{\partial}_b)(h_m^\lambda \otimes \omega_L^\lambda) = \sum_{|J|=q+1} \left( \sum_k \epsilon_{kL}^J d\pi_\lambda(\bar{Z}_k^\lambda) h_m^\lambda \right) \otimes \omega_J^\lambda,$$

and that

$$d\pi_{\lambda,\eta}(\bar{Z}_k^\lambda) = \sum_{|J|=q+1} \left( \sum_k \epsilon_{kL}^J (-\text{sgn } \mu_k) ((2m_k + 1)|\mu_k| - \mu_k)^{1/2} h_{m-(\text{sgn } \mu_k)e_k}^\lambda \right) \otimes \omega_J^\lambda.$$

Hence,

$$\begin{aligned} d\pi_\lambda(\bar{\partial}_b) \circ A_q^\lambda(h_m^\lambda \otimes \omega_L^\lambda) &= \\ (17) \quad &= \sum_{|J|=q+1} \sum_k \epsilon_{kL}^J \frac{(-\text{sgn } \mu_k)(|\mu_k|(2m_k + 1) - \mu_k)^{1/2}}{\alpha_L^\lambda + \sum_{j=1}^n (2m_j + 1)|\mu_j|} (h_{m-(\text{sgn } \mu_k)e_k}^\lambda \otimes \omega_J^\lambda) \end{aligned}$$

if  $m \neq 0$  or  $L \neq \bar{L}$  and  $d\pi_\lambda(\bar{\partial}_b) \circ A_q^{\lambda,\eta}(h_m^\lambda \otimes \omega_L^\lambda) = 0$  otherwise.

On the other hand,

$$d\pi_\lambda(\bar{\partial}_b)(h_m^\lambda \otimes \omega_L^\lambda) = \sum_{|J|=q+1} \sum_k \epsilon_{kL}^J (-\text{sgn } \mu_k) (|\mu_k|(2m_k + 1) - \mu_k)^{1/2} (h_{m-(\text{sgn } \mu_k)e_k}^\lambda \otimes \omega_J^\lambda),$$

so that

$$\begin{aligned} A_{q+1}^{\lambda,\eta} \circ d\pi_\lambda(\bar{\partial}_b)(h_m^\lambda \otimes \omega_L^\lambda) &= \\ (18) \quad &= \sum_{|J|=q+1} \sum_k \epsilon_{kL}^J \frac{(-\text{sgn } \mu_k)(|\mu_k|(2m_k + 1) - \mu_k)^{1/2}}{\alpha_J^\lambda - 2\mu_k + \sum_{j=1}^n (2m_j + 1)|\mu_j|} (h_{m-(\text{sgn } \mu_k)e_k}^\lambda \otimes \omega_J^\lambda). \end{aligned}$$

Notice that  $\lambda \in \Omega_q$  implies that  $\lambda \notin \Omega_{q+1}$ . If  $J = \{k\} \cup L$  as sets, then it is easy to check that  $\alpha_L^\lambda = \alpha_J^\lambda - 2\mu_k$ . When  $L = \bar{L}$  and  $k \notin \bar{L}$  then  $\mu_k > 0$  which implies that

$$d\pi_\lambda(K_{q+1}) \circ d\pi_\lambda(\bar{\partial}_b)(h_0^\lambda \otimes \omega_L^\lambda) = 0.$$

Thus, (17) and (18) prove equality (14) for  $\lambda \in \Omega_q$ .

The argument for  $\lambda \in \Omega_{q+1}$  is similar to the case  $\lambda \in \Omega_q$  and we omit it.  $\square$

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